New QCD Studies with the Resurrected JADE Data

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Outline

- Motivation
- The Experiment
- Revival of Data and Software
- QCD studies
 - Hadronic Event Shapes in e⁺e⁻ Annihilation
 - Strong Coupling Constant
 - Power Corrections
 - QCD Colour Factors
 - Longitudinal Cross Section
 - Momentum Spectra
- Conclusions

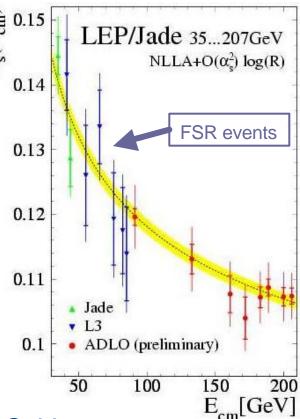
Motivation

- Explore perturbative and non-perturbative
 QCD effects at low energy scales Q
 - large leverage for predictions:

PT effects ∝ 1/log(Q)
NP effects ∝ 1/Q (typically

NP effects ∝ 1/Q (typically for event shapes)

interplay between hard and soft QCD best studied at "medium" energies



- JADE data: unique contribution for @ 14-44 GeV
- Test improved/new calculations from the LEP era at PETRA energies
 - New hadronic observables
 - New perturbative calculations
 - New MC models
 - New non-perturbative analytical approaches

as @ PETRA Times

1979 MARK-J Coll.:

– First direct measurement α_s based on LO for the Oblateness variable

1979+
$$a_s = 0.15 \dots 0.23$$
 @ $\sqrt{s} = 30 \text{ GeV}$

based on LO predictions

1982 CELLO Coll., JADE Coll.:

– first significant measurements of α_s NLO for Thrust and Differential 3 Jet Cross Section

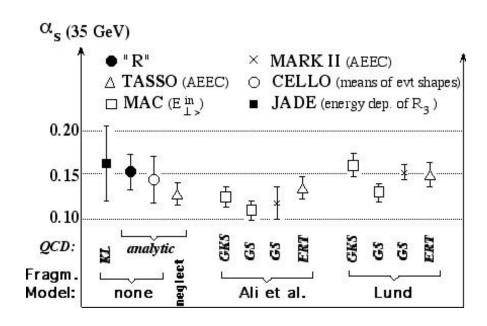
1982+
$$a_s$$
 (35GeV) = 0.11... 0.19

based on NLO predictions

...inconsistent results due to

- incomplete QCD matrix elements
- fragmentation models

Status of a_s in 1989



Summary value 1989:

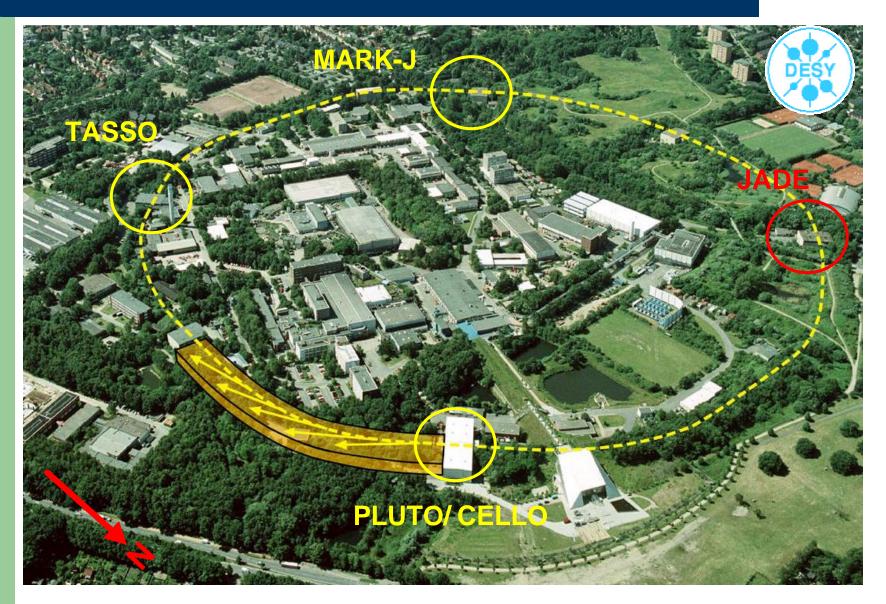
 $\alpha_{\rm S}$ (35GeV) = 0.14 ± 0.02

Use LEP techniques at PETRA energies

- to increase the precision
- to allow better comparison of results (values+systematics) over a wide range of e⁺e⁻ annihilation energies

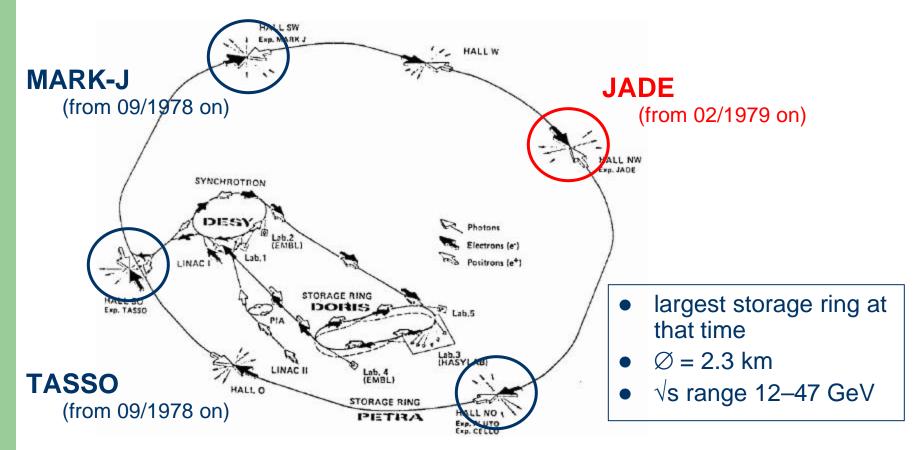
The Experiment

The PETRA e⁺e⁻ Storage Ring



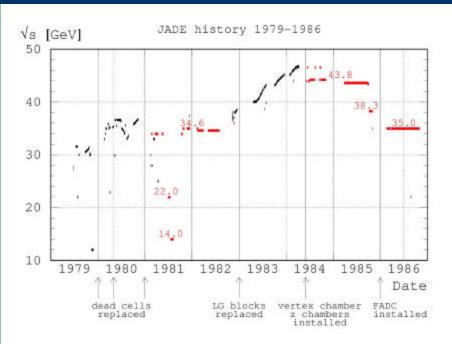
The PETRA e⁺e⁻ Storage Ring

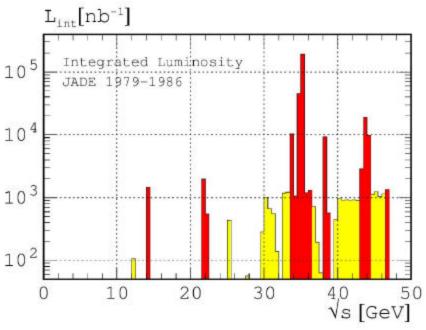
Operated 1978-1986 at DESY, Hamburg



PLUTO (from 09/1978 on)
CELLO (replacing PLUTO from 08/1982 on)

C.M.S. Energies and Luminosities





- Fixed energy runs
- Scan periods (Top quark search)
- By far most data accumulated at $\sqrt{s} = 35 \text{ GeV}$
- Total integrated lumi: 216 pb⁻¹
- Peak lumi: 24 μb⁻¹s⁻¹
 - \Rightarrow 26 multihadrons per hour @ σ^{had} =0.3 nb
- Clean multihadrons: 43100

The JADE Experiment

Participating states:

JApan (Tokyo), Deutschland (DESY, Hamburg, Heidelberg), England (Lancaster, Manchester, RAL), USA (Maryland)

≈ 120 collaborators in total

JADE is a magnetic, hermetic multipurpose detector:

Jet Chamber

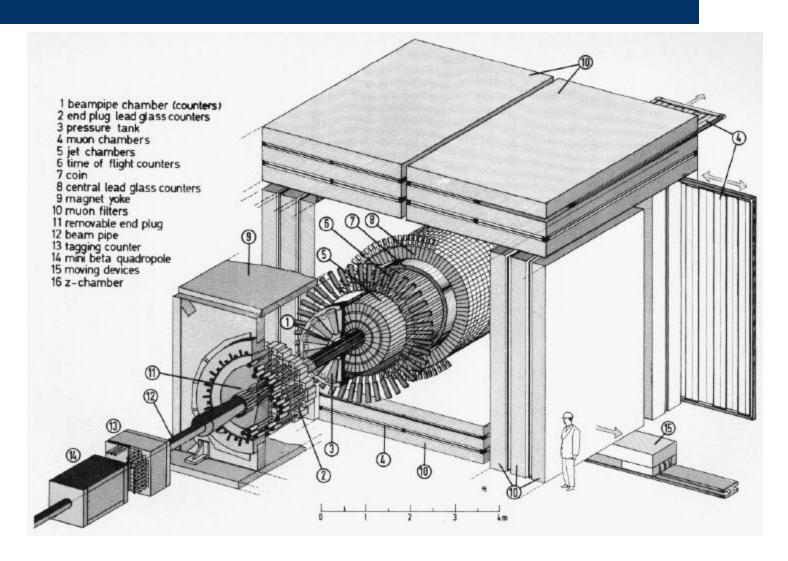
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Track curvature + dE/dx measurement, B = 0.48T 48 wire layers in r\varphi \sigma_{r\varphi}=180\mum (110\mum), \sigma_z=16-32mm (OPAL: \sigma_{r\varphi}=135\mum, \sigma_z=45-60mm)
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E.M. Calorimeter

```
\approx 2700 Lead Glass blocks (individually calibrated) 
 \sigma_{\rm E}/{\rm E} = 4%/\sqrt{\rm E}+1.5% 
 (OPAL: \sigma_{\rm E}/{\rm E} = 6.3%/\sqrt{\rm E}+0.2%)
```

Muon System
 up to 5 chamber layers / 3 absorber layers

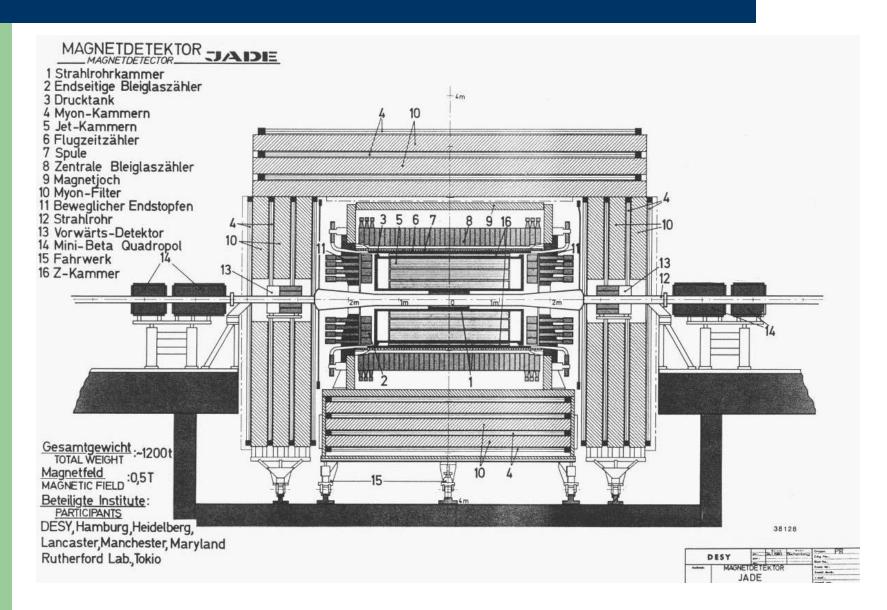
The Detector



Overall length/height: 8m/7m

(OPAL: 12m/12m)

The Detector



Revival of Data and Software

The JADE Revival Group

- RWTH Aachen, MPI Munich, DESY
 S. Bethke, O. Biebel, M. Blumenstengel, S. Kluth,
 P.A.M.F., C. Pahl, P. Pfeifenschneider,
 and J.E. Olsson
- Since 1998: 20+ publications/conference contributions based on/involving the reanalysed JADE data
- New JADE results considered in numerous publications from LEP collaborations / QCD theory groups
- Inspires a LEP working group to address the difficult question of keeping data and software of LEP collaborations alive

Resurrection of the JADE Data ...

- Original data were located on
 - IBM mainframe at the DESY computer centre
 - IBM tapes at DESY/Heidelberg U.
- DESY IBM completely closed July 1997
 - Last-Minute transfer to "modern" data carriers (IBM/EXABYTE cartridges) and computer platforms
- Now: data partially reside on CERN Castor tapes, DVDs ...
- Data organisation mainly based on the data management system BOS (version 1979)
 - Raw Data (REDUC1/REDUC2): BOS banks converted into FPACK (platform independent, still need to reconvert)
 - MH data sets (ZE4V) converted into ASCII (used for reanalyses)

... and of the JADE Software

- Detector simulation
 - detailed particle tracking, detector response, inefficiencies, resolution
- Event analysis software
 - pattern recognition, cluster analysis ...
- JADE interactive graphics
 - event display, event analysis, event editing
- MH filtering and packing software

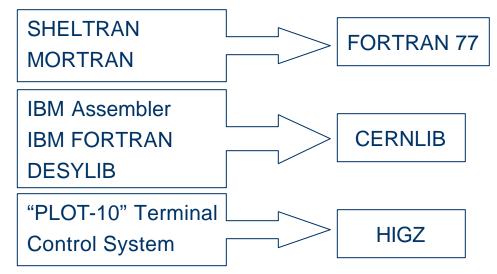
Source Code

- Code fragments date from 1974 on
- Mixture of different FORTRAN standards (FORTRAN IV, FORTRAN 77)
- "Illegal" IBM extensions
- Ancient pre-compiler languages (SHELTRAN, MORTRAN)
- IBM/370 assembler code

Big parts are extremely unstructured "spaghetti" code, badly documented!

Tasks

- Extract knowledge and information from incoherently spread sources (nontrivial "archaeological" challenge)
 Code modification
 >JADE Computer Notes
 JADE Notes
 JADE PhD theses
 Manual fragments
 Source code
- Emulation interfaces
 (missing libraries, IBM FORTRAN intrinsics ...)

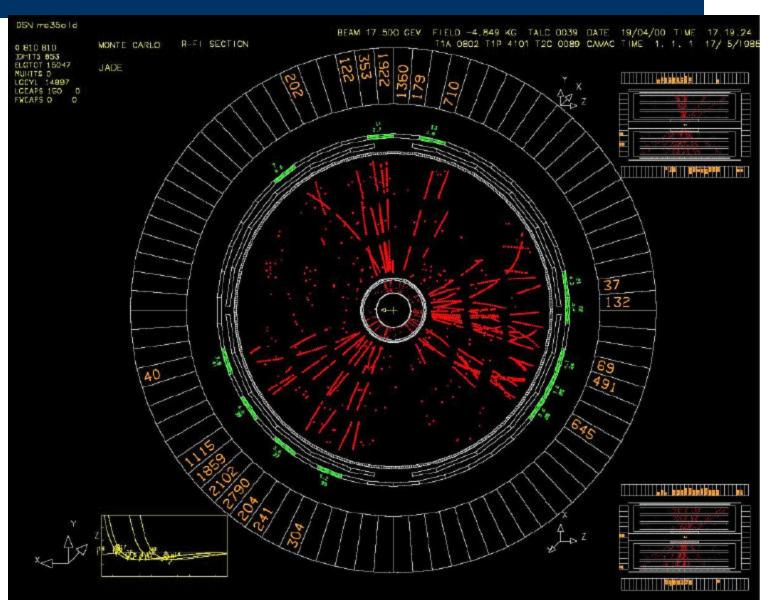


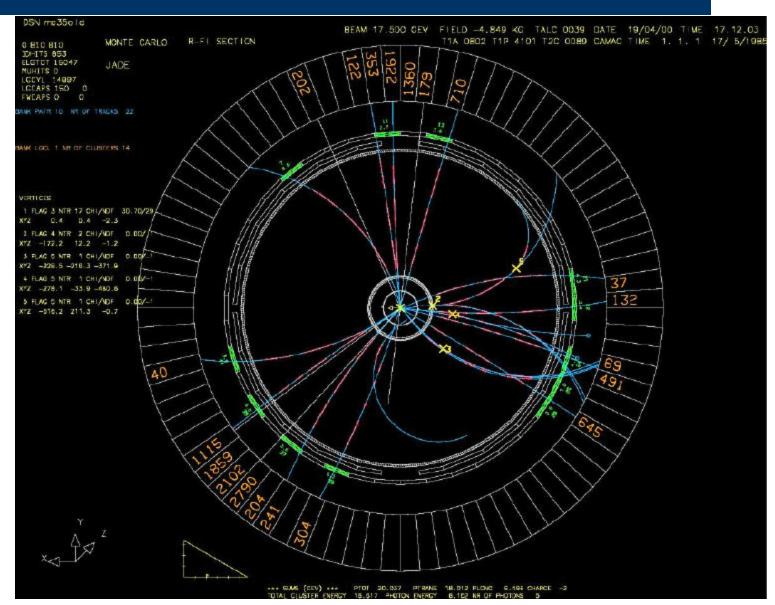
Platform dependent features extremely problematic!!!

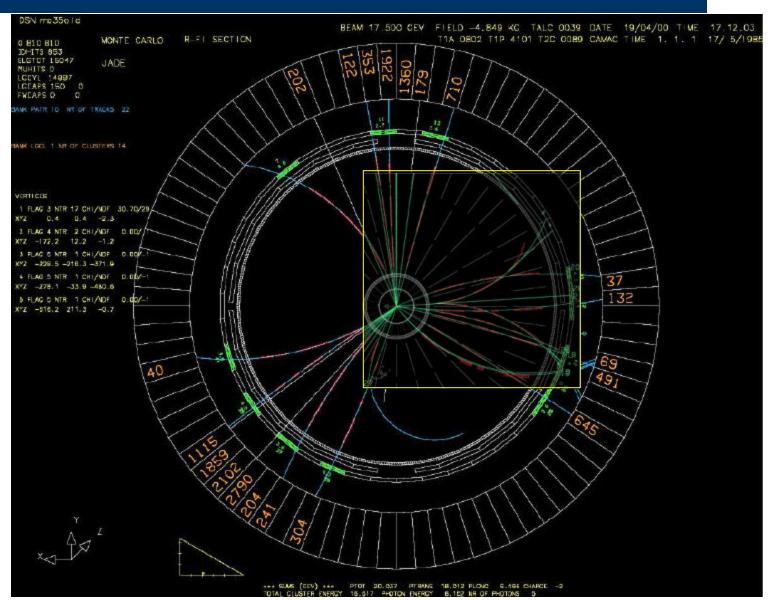
- Bit&Byte manipulation
- Endian convention (byte storage order)

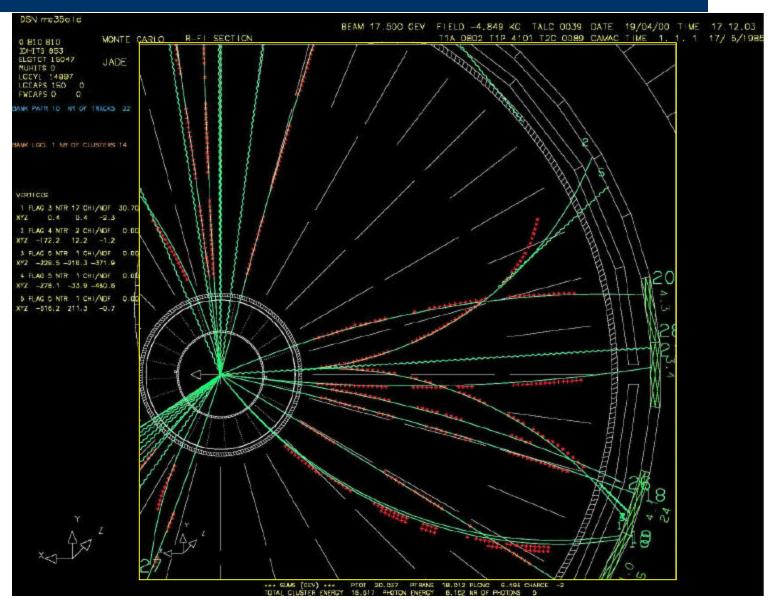
Complete installation succeeded on IBM RS/6000 AIX!

- XLF compiler advantageous
- same endian scheme as IBM/370









Performance I (Jet Chamber)

Pythia Jetset(J) Ariadne

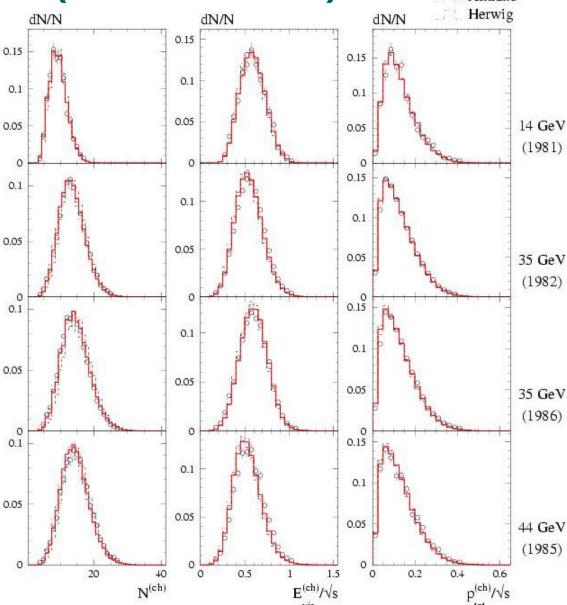
♦ JADE

Integral quantities:

N^(ch),

 $E_{vis}^{(ch)}/\sqrt{s}$, $p_{tot}^{(ch)}/\sqrt{s}$,

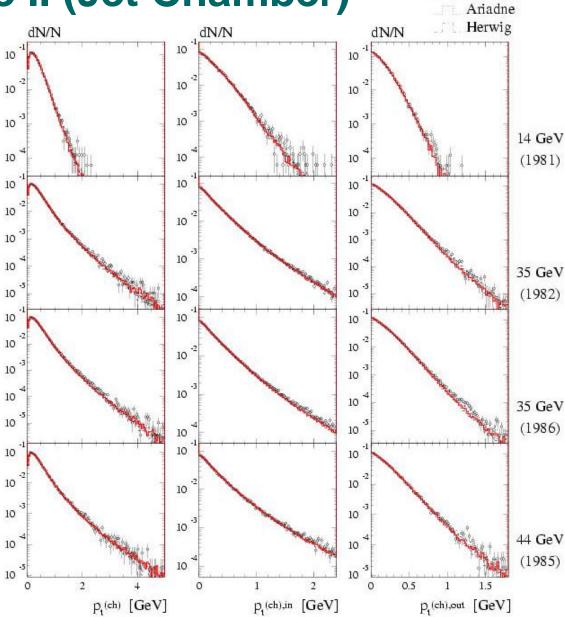
MPI Colloquium "QCD Studies with the Resurrected JADE Data"



Performance II (Jet Chamber)

Particle spectra:

 $p_t^{(ch)}$, $p_t^{(ch), in}$, $p_t^{(ch), out}$,



June 17th, 2003

Jetset(J)

Performance III (Lead Glass)

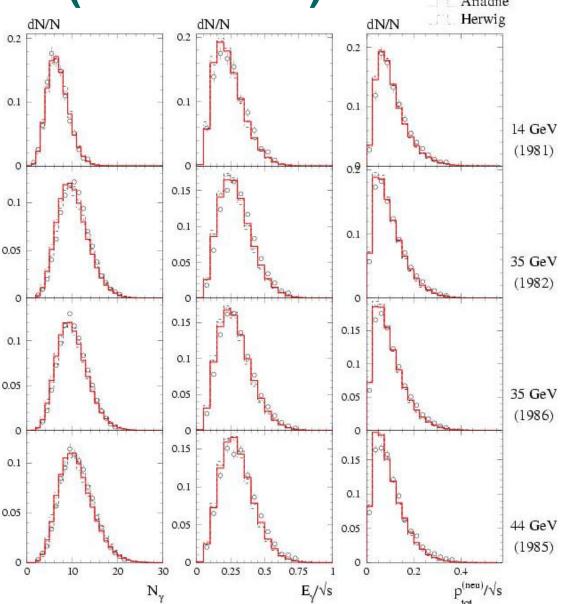
L Pythia Jetset(J) Ariadne

♦ JADE

Integral quantities:

 N_{γ} , E_{γ}/\sqrt{s} , $p_{tot}^{(neu)}/\sqrt{s}$,

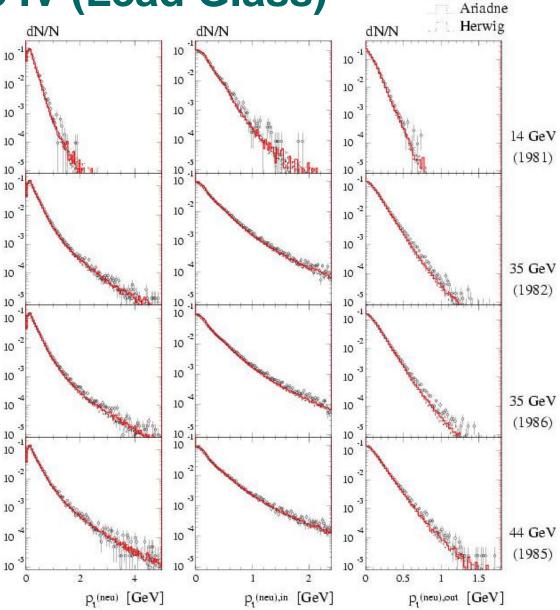
• • •



Performance IV (Lead Glass)

Particle spectra:

p_t(neu), p_t(neu), in, p, (neu), out,



♦ JADE Pythia

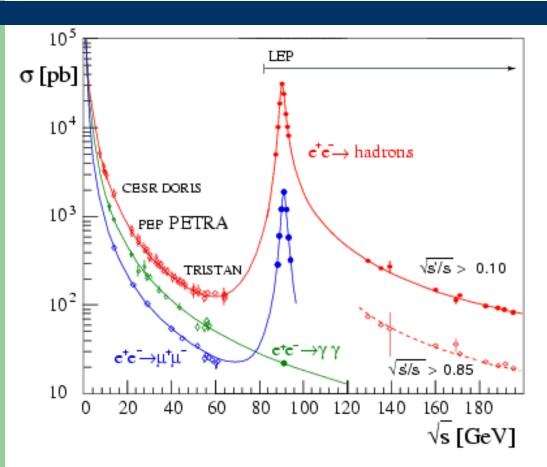
Jetset(J)

Revival Summary

- JADE software works reliably
- JADE simulation capable of reproducing most integral observables and particle spectra measured with the real detector
- JADE simulation usable for the correction of physical quantities, e.g.:
 - Event shape observables
 - Momentum spectra
 - ...

QCD Studies

Hadronic Final States

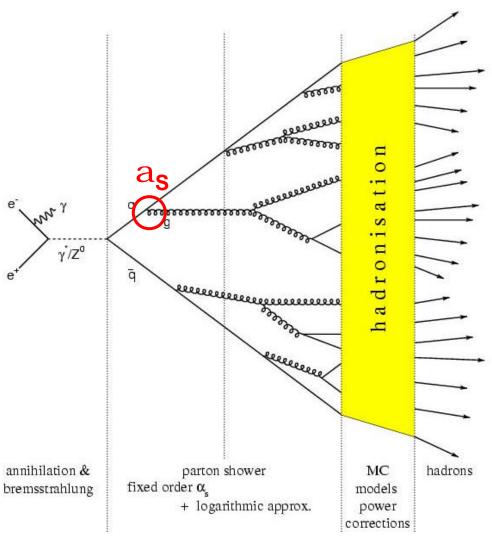


Cross section for e⁺e⁻ → hadrons:

- σ^{had} (PETRA) = 0.1...10nb $\approx 1/100\sigma^{had}$ (M_Z)
- Hadron production at PETRA energies mainly via γ^* exchange

June 17th, 2003

QCD in e⁺e⁻ Annihilation



PT QCD:

- $O(\alpha_S^2)$, NLLA, ...
- Parton shower MC

NP QCD:

- Phenomenological hadronisation models
- Analytical power corrections

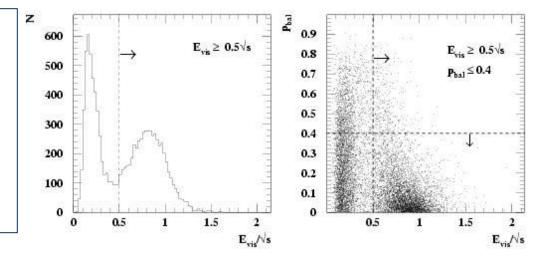
35 GeV

1 GeV

Multihadronic Selection

Main Selection Cuts:

- 4 tracks from vertex region
- 3 "long + good" tracks
- Visible Energy > 0.5.√s
- Momentum balance < 40%
- Missing Momentum < 0.3√s
- $|\cos \Theta_{\mathsf{T}}| < 0.8$



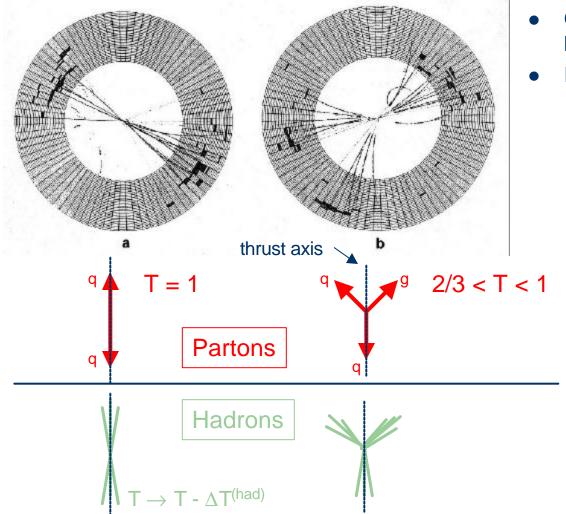
Residual background ≈ 1%

- $e^+e^- \rightarrow e^+e^- \gamma \gamma$
- $e^+e^- \rightarrow \tau^+\tau^-$

MH data samples for main analyses:

\sqrt{s} -range [GeV]	data taking period	\mathcal{L} $[pb^{-1}]$	$\langle \sqrt{s} \rangle$ [GeV]	MH data
14.0	JulAug. 1981	1.46	14.0	1734
22.0	JunJul. 1981	2.41	22.0	1390
33.8 - 36.0	Feb. 1981 - Aug. 1982	61.7	34.6	14372
35.0	FebNov. 1986	92.3	35.0	20925
38.3	OctNov. 1981	8.28	38.3	1587
43.4 - 46.6	Jun. 1984 - Oct. 1985	28.8	43.8	3940

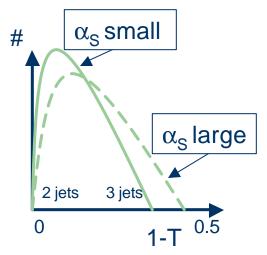
Hadronic Event Shapes



- Quantify the shape of an event by a single number.
- Example: "Thrust"

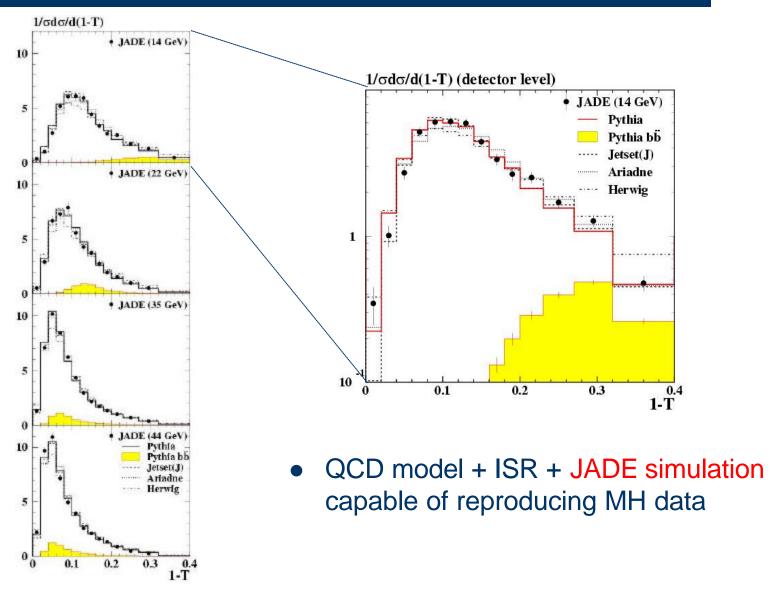
$$T = \max_{\vec{n}} \left(\frac{\sum_i |\vec{p_i} \vec{n}|}{\sum_i |\vec{p_i}|} \right)$$

QCD expectation:



Event shape observables are sensitive to PT and NP effects!

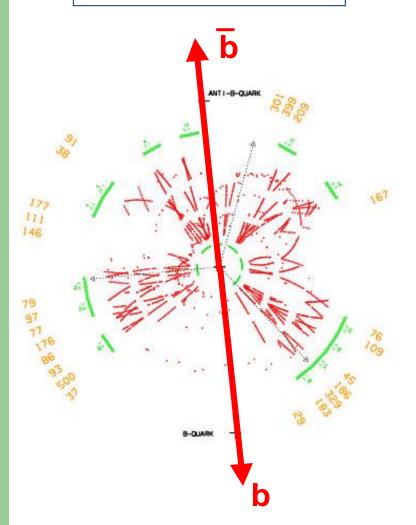
Detector Level Distributions



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bb Events

Pythia event @ 14 GeV



- 9% fraction
- fake hard gluon radiation due to electroweak decays + mass effects
- 14 GeV: up to 50%
 contamination in extreme
 3 jet region

Treat as "background" in view of later comparison with massless QCD calculations!

More Event Shapes

Thrust T

Thrust
$$T$$

$$T = \max_{\vec{n}} \left(\frac{\sum_{i} |\vec{p_i} \vec{n}|}{\sum_{i} |\vec{p_i}|} \right) \Rightarrow \text{thrust axis } \vec{n_T} \text{ event hemispheres } H_k \Rightarrow B_k = \frac{\sum_{i \in H_k} |\vec{p_i} \times \vec{n_T}|}{2\sum_{i} |\vec{p_i}|}, k = 1, 2$$

Heavy Jet Mass M_H

$$M_{\rm H}^2 = \frac{\max(M_1^2, M_2^2)}{(\sum_i E_i)^2}$$

Total/Wide Jet Broadening B_T , B_W

$$B_{\text{T}} = B_1 + B_2$$

 $B_{\text{W}} = \max(B_1, B_2)$

C Parameter

$$\begin{split} \Theta^{\alpha\beta} &= \frac{\sum_i (p_i^\alpha p_i^\beta)/\left|\vec{p_i}\right|}{\sum_i \left|\vec{p_i}\right|}\;, \quad \alpha,\,\beta = 1,\,2,\,3 \\ &C = 3(\lambda_1\lambda_2 + \lambda_2\lambda_3 + \lambda_3\lambda_1) \end{split}$$

• Calculate eigenvalues λ_i from linearised momentum tensor.

Differential 2 Jet Rate y_{23} (Durham Scheme)

$$y_{ij} = \frac{2 \min(E_i^2, E_j^2)(1 - \cos \vartheta_{ij})}{(\sum_k E_k)^2}$$

 $\frac{dR_2(y_{\text{cut}})}{dy_{\text{cut}}} = \frac{1}{\sigma} \frac{d\sigma(y_{23})}{dy_{23}}$.

- Define jet resolution parameter y_{ii}.
- Combine particles i, j with smallest y_{ij} into pseudo particles and proceed until $y_{ij} > y_{cut}$ for 2 remaining pseudo particles ("jets").

Measurement

Observables: y=1-T, M_H , B_T , B_W , C, y_{23}

- Infrared and collinear safe quantities
- Resumable in all orders $\alpha_s \log (1/y)$ (important in 2 jet region)

Perform MC based corrections to measured distributions

- bb-fraction on detector level
 - reduces mass effects
- Detector effects
 - Resolution, acceptance, secondary processes
- MH selection
 - acceptance
- Photon ISR

Hadron level distributions comparable with QCD predictions

QCD Models

- PYTHIA/JETSET:
 - LLA parton shower + string fragmentation
- ARIADNE:
 - colour dipole scheme + string fragmentation
- HERWIG:
 - MLLA parton shower + cluster fragmentation
- COJETS:
 - LLA parton shower + independent fragmentation

Use LEP versions tuned to OPAL data

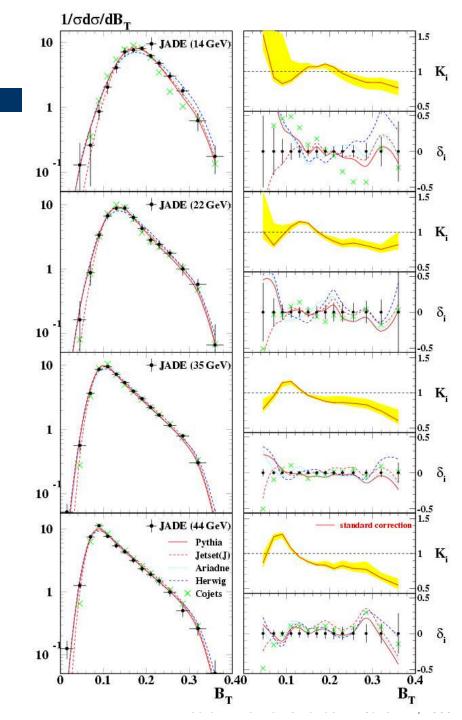
Try also former JADE optimisation for JETSET 6.3

Hadron Level

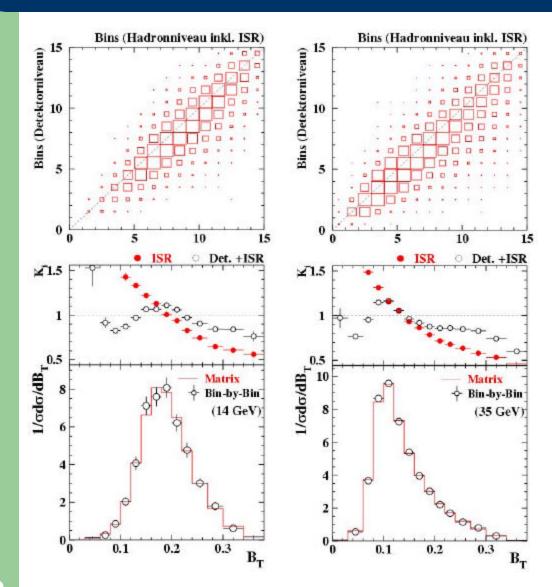
Bin-by-bin unfolding with correction factors $K_i = MC_i^{had}/MC_i^{det}$ based on udsc samples:

- PYTHIA
 - good overall consistency
- HERWIG/ARIADNE
 - moderate at 14+22 GeV, better at higher √s
- JETSET (JADE)
 - good at 14+22 GeV, slightly worse at higher √s
- COJETS
 - disfavoured at 14+22 GeV, remains worse at higher √s

Event shape become more and more 2 jet like at higher energies



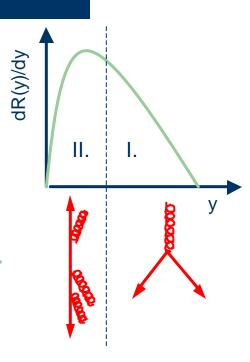
Matrix vs. Bin-by-Bin Unfolding



- Consistent hadron levels
- Detector effects partially compensate ISR

Determination of the a_s

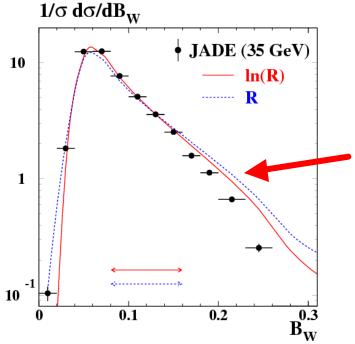
- PT prediction for the cumulative cross section $R(y) = \int_{0}^{y} dy' 1/\sigma \cdot d\sigma/dy'$
- I. NLO: describes "hard" gluon contribution $R(y)=1+A(y)\cdot\alpha_s+B(y)\cdot\alpha_s^2$
- II. NLLA: describes "soft" gluon contribution $R(y) = (1 + C_1 \cdot \alpha_S + C_2 \cdot \alpha_S^2) \exp\{Lg_1(\alpha_S L) + g_2(\alpha_S L)\}$ L = In(1/y)
- III. Combination of NLO+NLLA, e.g.: In(R) matching In(R) = Lg₁(α_s L)+g₂(α_s L) $-(G_{11}$ L+G₁₂L²)· α_s -(G₂₂L+G₂₃L²)· α_s ²
 +A(y)· α_s +[B(y)- ½ A(y)²]· α_s ²

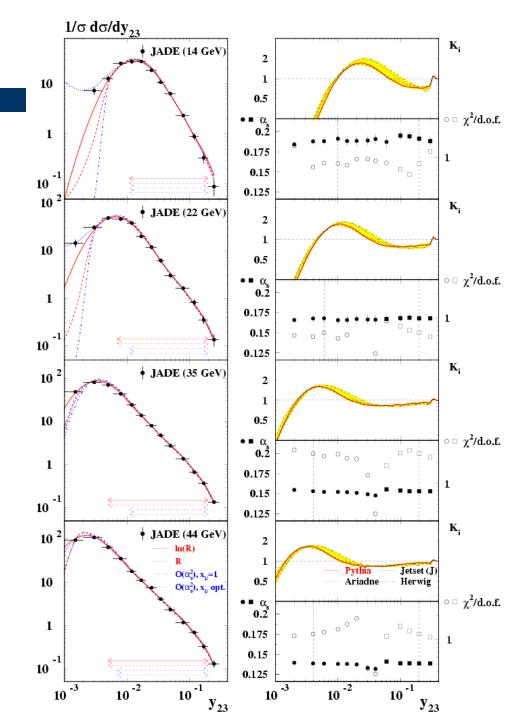


- NP effects: PYTHIA, JETSET(JADE), ARIADNE, HERWIG
- Fit α_s with renormalisation scale factor $x_{\mu}=\mu/\sqrt{s}=1$ + bin-by-bin hadronisation correction of R(y) (standard=PYTHIA)

Fit Curves

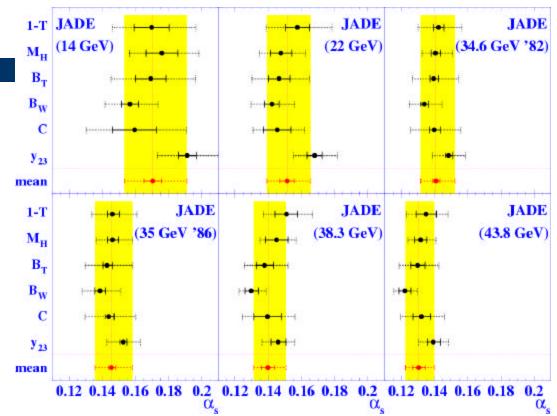
- Typically $\chi^2/d.o.f. = 0.5...2.0$
- Stable Fits
- Large hadronisation corrections at 14 GeV!
- Problems with B_W





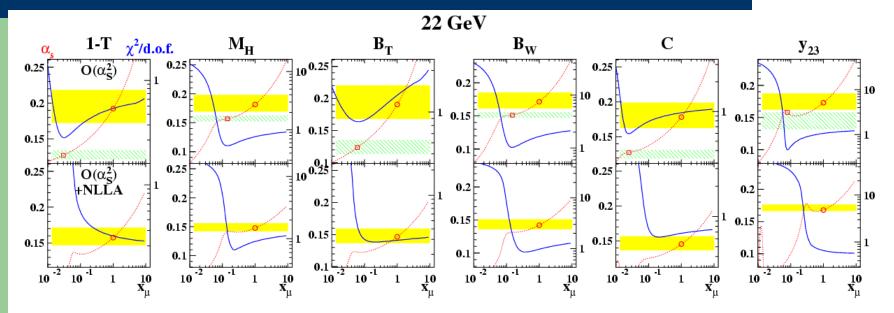
a_s Results

- Similar scattering of individual results due to missing higher order terms, but...
- ...results consistent within
 1-2σ of experimental errors
- x_μ dependence significantly smaller w.r.t. pure NLO results!
- Dominant errors:
 - Renormalisation scale
 - Hadronisation (14+22GeV!)
 - Mass effects (14+22GeV!)



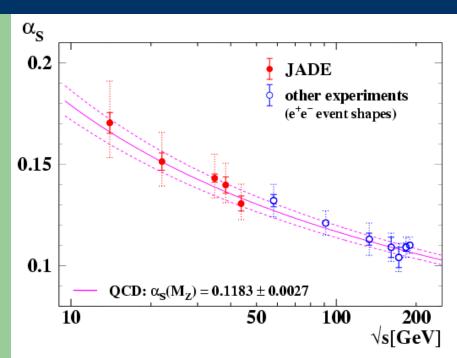
$\langle \sqrt{s} \rangle$ [GeV	$]$ $\alpha_{\rm S}(\sqrt{s})$	fit error	exp.	hadr.	higher ord.	total
14.0	0.1704	±0.0	051*	$^{+0.0141}_{-0.0136}$	$^{+0.0143}_{-0.0091}$	$^{+0.0206}_{-0.0171}$
22.0	0.1513	±0.0	043*	± 0.0101	$^{+0.0101}_{-0.0065}$	$^{+0.0144}_{-0.0121}$
34.6 ('82)	0.1409	±0.0012	± 0.0017	± 0.0071	$^{+0.0086}_{-0.0057}$	$^{+0.0114}_{-0.0093}$
35.0 ('86)	0.1457	±0.0011	±0.0020	±0.0076	$^{+0.0096}_{-0.0064}$	$^{+0.0125}_{-0.0101}$
38.3	0.1397	± 0.0031	±0.0026	± 0.0054	$^{+0.0084}_{-0.0056}$	$^{+0.0108}_{-0.0087}$
43.8	0.1306	±0.0019	±0.0032	± 0.0056	$^{+0.0068}_{-0.0044}$	$^{+0.0096}_{-0.0080}$

Renormalisation Scale



- NLO+NLLA: reduced x_{μ} dependence around x_{μ} =1 compared to NLO
 - $\alpha_{\rm S}(\sqrt{\rm s}, {\rm x_u}=1)$ more consistent than in NLO case
 - But: sizable α_S dependence around x_u =1 still present
- Pure NLO: Preference for small $x_{\mu}^{(opt)} = O(0.01...0.5)$
 - scale dependence around $x_{\mu}^{(opt)}$ sometimes smaller, but...
 - less consistent individual results
 - (α_S, x_u) fits not always stable, large statistical errors
 - no strong theoretical arguments for the choice $x_{\mu} = x_{\mu}^{(opt)}$
 - \Rightarrow have to consider **both** $\alpha_{S}(\sqrt{s}, x_{\mu}=x_{\mu}^{(opt)})$ **and** $\alpha_{S}(\sqrt{s}, x_{\mu}=1)$ NLO+NLLA @ $x_{\mu}=1$ seems to be the "natural" choice

Test of the Running of as



 QCD fit, exp.+stat. uncertainties (inner error bars):

$$\Lambda_{MS}^{(5)} = 246 \pm 7 \text{ MeV}$$

 $\alpha_{S}(M_{Z}) = 0.1210 \pm 0.0006$
 $P(\chi^{2}) = 75\%$

 α_S = const., total errors (outer error bars):

$$P(\chi^2) = 1.1 \cdot 10^{-5}$$

- Now more values with higher accuracy available
- α_S of "homogeneously" determined from PETRA to LEP2 energies

QCD expectation:

$$C_{\Delta}=3$$
, $C_{F}=4/3$, $N_{F}=5$

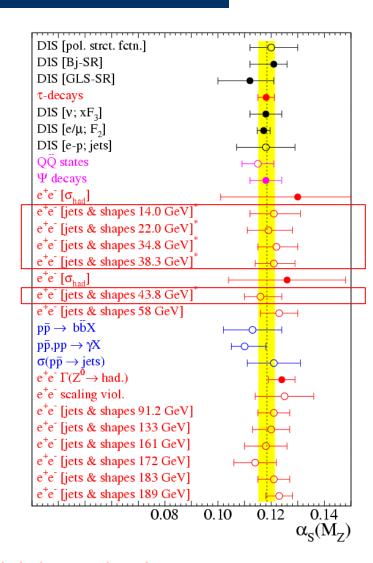
$$\begin{split} \alpha_{\rm S}(\sqrt{s}) &= \frac{1}{\beta_0 l} - \frac{\beta_1 \ln l}{\beta_0^3 l^2} + \frac{1}{\beta_0^3 l^3} \left[\frac{\beta_1^2}{\beta_0^2} (\ln^2 l - \ln l - 1) + \frac{\beta_2}{\beta_0} \right] \\ l &= \ln(\sqrt{s}/\Lambda_{\overline{\rm MS}})^2 \\ \beta_0 &= \frac{1}{12\pi} \left(33 - 2N_f \right) \\ \beta_1 &= \frac{1}{24\pi^2} \left(153 - 19N_f \right) \\ \beta_2 &= \frac{1}{3456\pi^3} \left(77139 - 15099N_f + 325N_f^2 \right) \end{split}$$

Good agreement with world average based on NNLO QCD

as Summary

- LEP established resummed calc for event shape work well at PETRA energies
- LEP tuned MC models (PYTHIA) capable of describing data down to 14 GeV
- Consistent picture of individual α_s results
- Hadronisation uncertainties at 14 GeV as large as renormalisation scale ambiguity
- New PETRA results now better comparable with LEP (values+systematics)
- Results consistent with other measurements and methods

```
\alpha_{\rm S}(M_{\rm Z^0}) = 0.1194^{+0.0083}_{-0.0070}
                                              (PETRA)
\alpha_{\rm S}(M_{\rm Z^0}) = 0.121 \pm 0.006 \; ({\rm LEP + SLC})
\alpha_{\rm S}(M_{\rm Z^0}) = 0.120 \pm 0.007 \text{ (LEP2)}
```



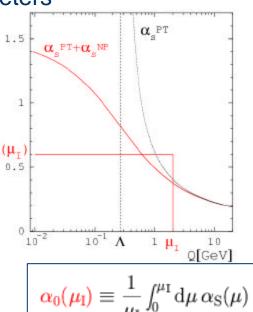
Power Corrections

- Classical method to estimate NP effects: MC models
 - PYTHIA, HERWIG, ARIADNE ...
 - numerous parton shower + fragmentation parameters
- Promising alternative: "power corrections"
 - Parametrise unknown but analytical behaviour of the physical strong coupling constant around the Landau pole Λ (0...2GeV)
 - Dokshitzer, Marchesini, Webber (DMW): NP structure due to soft gluon radiation at $\mu \approx \Lambda$

$$\langle y \rangle = \langle y \rangle^{\rm PT} + \mathcal{D}_y \mathcal{P}$$
 (means)
$$\frac{\mathrm{d}\sigma}{\mathrm{d}y}(y) = \frac{\mathrm{d}\sigma^{\rm PT}}{\mathrm{d}y}(y - \mathcal{D}_y \mathcal{P})$$
 (distributions)

$$\mathcal{P} = \frac{4C_F}{\pi^2} \mathcal{M} \frac{\mu_{\rm I}}{Q} \left[\alpha_0(\mu_{\rm I}) - \alpha_{\rm S}(\mu_{\rm R}) - \beta_0 \frac{\alpha_{\rm S}^2(\mu_{\rm R})}{2\pi} \left(\ln \frac{\mu_{\rm R}}{\mu_{\rm I}} + \frac{K}{\beta_0} + 1 \right) \right]$$

- α_0 is the only NP parameter!
- α_0 is universal



$$\frac{\alpha_0(\mu_{\rm I})}{\alpha_0} \equiv \frac{1}{\mu_{\rm I}} \int_0^{\mu_{\rm I}} \mathrm{d}\mu \, \alpha_{\rm S}(\mu)$$

Power Corrections to Distributions

Observable specific part is D_v:

 \Rightarrow T, M_H, C: shift

 \Rightarrow B_T, B_W: shift+squeeze

(y₂₃: no 1/Q contribution)

y	$\mathcal{D}_y = \mathcal{D}_y(\alpha_{\mathrm{S}}, y)$
1-T	2
$M_{ m H}^2$	1
C	3π
B_{T}	$\ln(1/y) + D_T(y, \alpha_S(yQ))$
B_{W}	$\frac{1}{2}\ln(1/y) + D_1(y, \alpha_{\mathrm{S}}(yQ))$

Test of DMW ansatz:

- Use mod. ln(R) matching for PT part
- Perform simultaneous (α_s, α_0) fits to all available event shape spectra Available data sets:

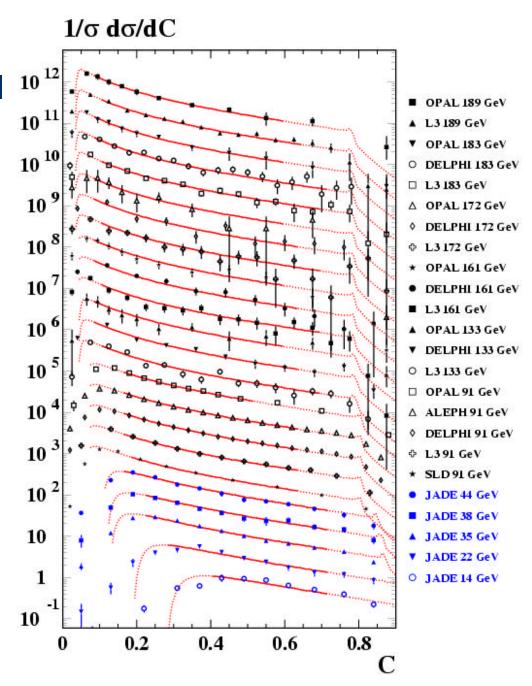
Accelerator	\sqrt{s} [GeV]	1-T	$M_{\rm H}$ $B_{\rm T}$, $B_{\rm W}$, C
PETRA (JADE, TASSO)	12-47	10200	0 43700
PEP (HRS, MARK II)	29	28300	
TRISTAN (AMY)	55-58	1900	
LEP I (ADLO*)	91		$O(10^6)$
SLC (SLD)	91		37200
LEP II (ADLO*)	133-189		15600

JADE is the only contribution for new observables below M₇

... covering the energy range $\sqrt{s} = 14...189$ GeV!!!

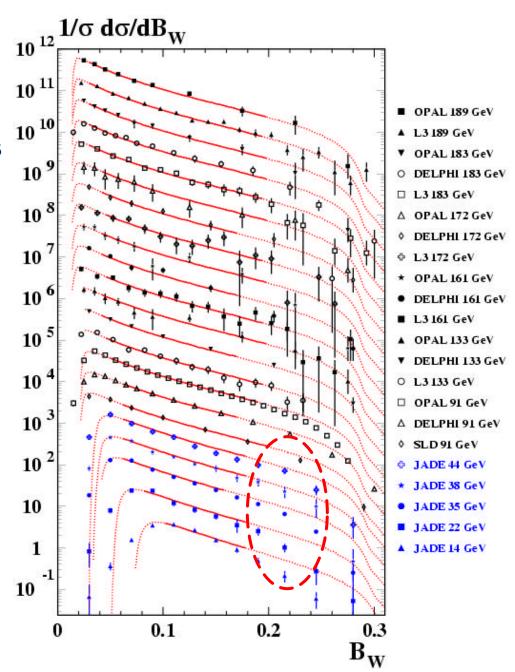
DMW Fits (I)

Good description of the data (T, C, B_T) within the kinematical limit of the predictions

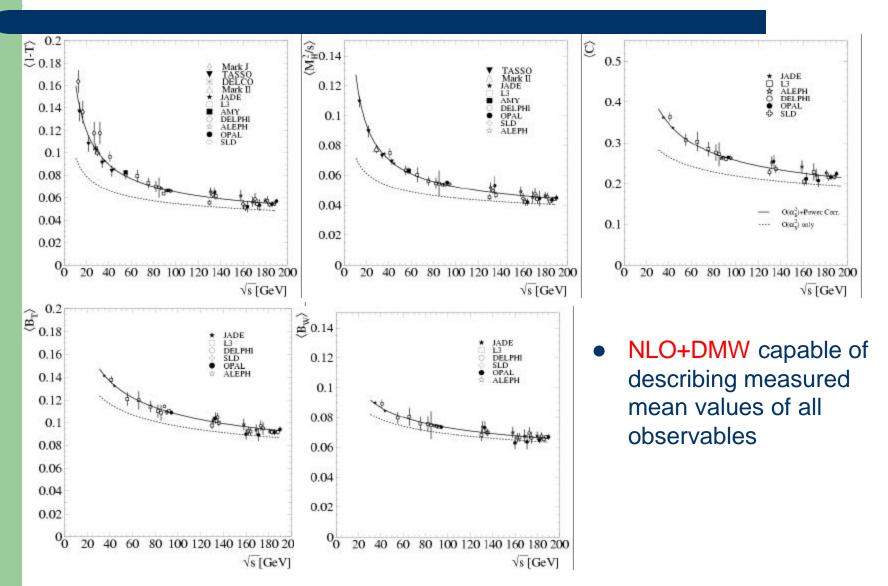


DMW Fits (II)

- Excess in 3 jet region for less inclusive observables (M_H,B_W) at PETRA energies!
- NB: also problems with PT prediction for B_W

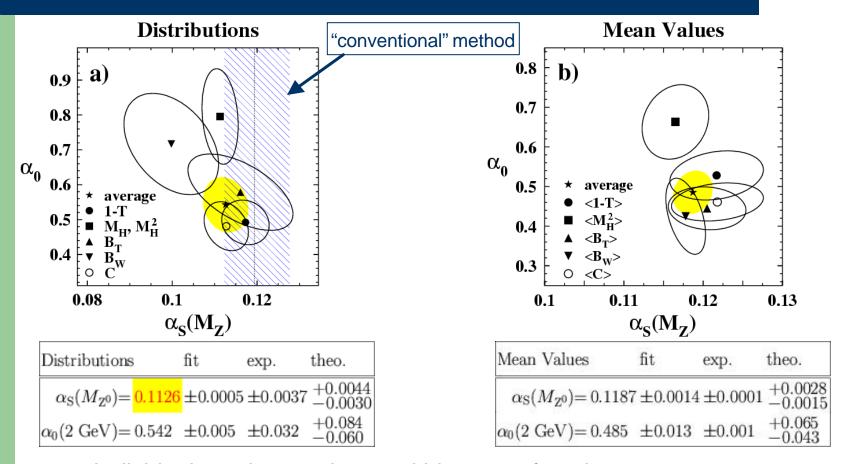


DMW Fits to Mean Values



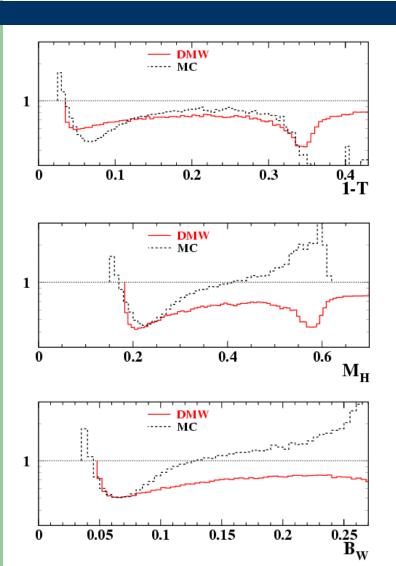
[Does not include update at 14+22 GeV]

(a_s,a_0) -Results



- Individual results consistent within 1-2σ of total errors
- α_0 universal within 20% uncertainty level of the Milan factor (stemming from $O(\alpha_s^2)$ evaluation of power corrections)
- But: $\alpha_s^{(pow.corr)} < \alpha_s^{(MC)}$ due to minor/missing squeeze of PT spectrum (fit chooses small α_s to compensate; big effect for jet broadening variables!!!)

Power Corrections vs. MC Predictions

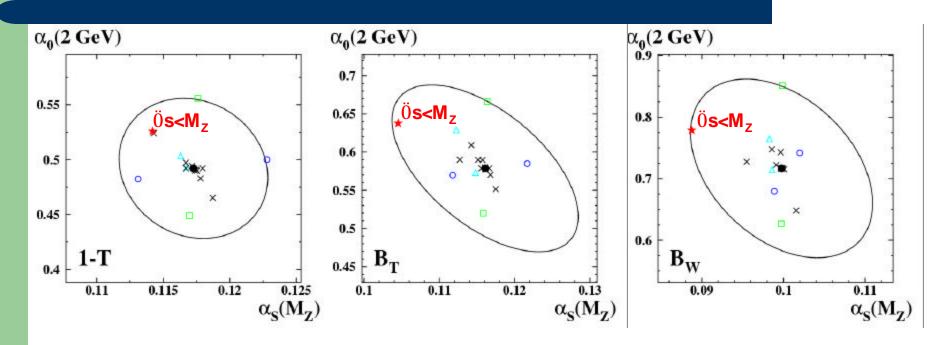


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PC/MC corrections expressed by means of corrections factors:

- T, C, B_T with "similar" corrections
- M_H, B_W with strongly deviating corrections
- "Missing squeeze" (w.r.t. MC prediction) is compensated by small α_s values

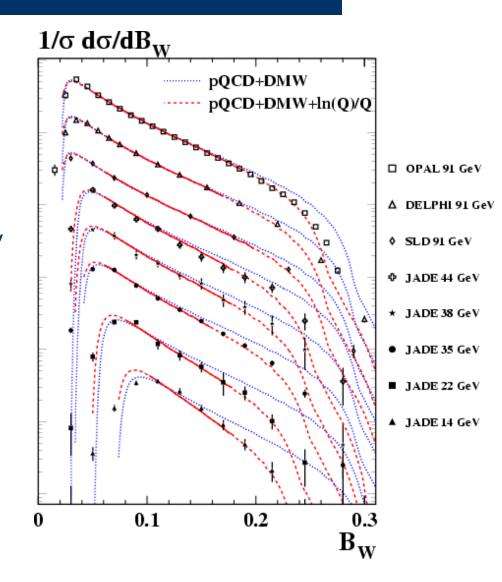
Missing NP Terms?



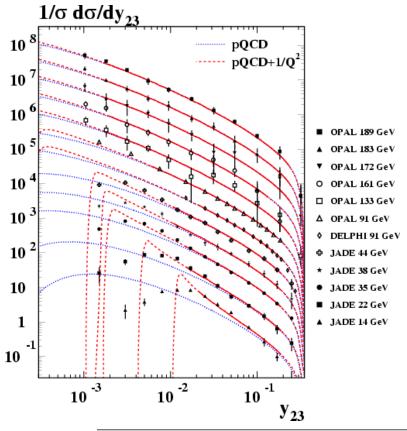
- Explore possible missing (higher order) terms by fits to separate data sets: √s<M_z, √s ³M_z, ...
 - Large systematic effects for B_T, B_W

Extended Power Corrections

- Evidence for additional terms probably behaving ∞ ln(Q)/Q
 - Extended power corrections?
 - Missing PT terms?
 (effect partially reproduced by redefining x_μ)
- Log enhanced power corrections expected due to mass effects (but expected effect for B_w not as large)



Power Corrections to y₂₃



- DMW: 1/Q coefficient = 0
 ...confirmed by fit
- Evidence for additional terms probably behaving

 ¹/Q²
- Need 14+22 GeV data to see the effect!

		$\alpha_S(M_{Z^0})$	$A_{10}[GeV]$	$A_{20}[\text{GeV}^2]$	$\chi^2/\mathrm{d.o.f.}$
I	pQCD	0.1147 ± 0.0005	_	_	59.7/100
	pQCD	0.1152 ± 0.0005	_	_	151/107
П	$pQCD+A_{10}/Q$	0.1124 ± 0.0006	0.062 ± 0.008	_	98.2/106
	$pQCD+A_{20}/Q^2$	0.1133 ± 0.0005		2.25 ± 0.18	71.2/106
	$pQCD+A_{10}/Q + A_{20}/Q^2$	$0.1128{\pm}0.0007$	0.018 ± 0.014	1.94 ± 0.31	69.7/105

Power Corrections Summary

- PETRA data discriminate between "good" (T, C, B_T) and "bad" (M_H, B_W) observables (w.r.t. of DMW model)
- α_0 universal at 20% level
- DMW (for distributions) different from MC prediction

$$\rightarrow \alpha_{\text{S}}^{\text{(pow.corr)}} < \alpha_{\text{S}}^{\text{(MC)}}$$

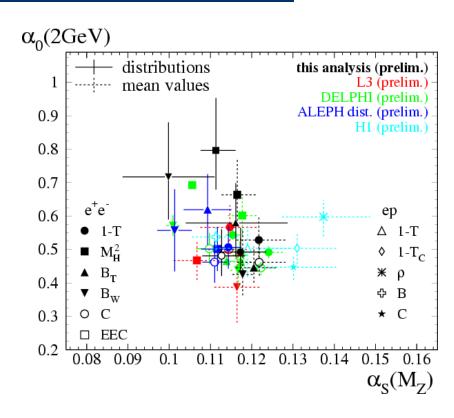
- Indication of higher order terms
 (B_W, y₂₃) may inspire theorists?
- Combined means+distributions:

$$\alpha_{\rm S}(M_{\rm Z^0}) = 0.1175^{+0.0031}_{-0.0021}$$

 $\alpha_{\rm 0}(2~{\rm GeV}) = 0.503^{+0.066}_{-0.045}$

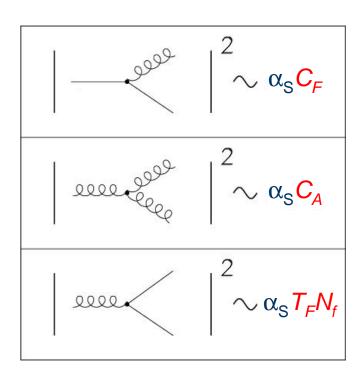
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Consistency with other measurements



More and improved PC calculations needed!

Colour Factors from Event Shapes



Relative weights of fundamental vertices determined by QCD gauge structure:

$$C_F = 4/3$$
, $C_A = 3$, $T_f N_f = 1/2N_f$

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Colour structure known for event shape

PT part

$$A \propto C_F$$
, $B = B(C_A, C_F, N_F)$
 $NLLA = NLLA (C_A, C_F, N_F)$

- Running α_S $\beta_0 = \beta_0 (C_A, N_F), \beta_1 = \beta_1 (C_A, C_F, N_F)$
- Power Corrections

$$P = P(C_A, C_F, N_F)$$

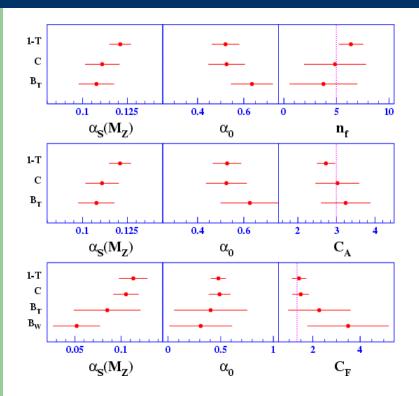
$$M = M(C_A, N_F)$$

$$D_V = D_V(C_A, C_F, N_F)$$

Reduced model dependence!

(i.e. no bias from colour structure of MC)

Results



	Fit α_S and α_0 and		Fix α_0 and N_f and		
	$(C_A \text{ or } C_F \text{ or } N_f)$		fit $\alpha_{\rm S}$ and C_A and C_F		
	1-T	C	1-T	C	QCD
			2.7 ± 0.2	3.0 ± 0.5	3
C_F	1.4 ± 0.3	1.5 ± 0.4	1.3 ± 0.2	1.3 ± 0.5	4/3
N_f	6.4 ± 1.2	4.9 ± 3.0	_	_	5

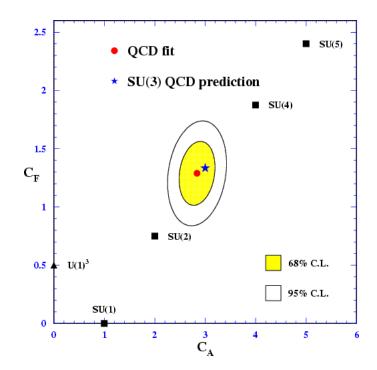
Combined results:

$$C_F = 2.84 \pm 0.24$$

$$C_A = 1.29 \pm 0.18$$

...competitive with 4 jet angular correlation analyses

Need JADE data to constrain the fit



Longitudinal Cross Section S₁

Differential cross section for inclusive hadron production in e+e-® g,Z ® h+X

$$\frac{1}{\sigma_{\rm tot}} \cdot \frac{\mathrm{d}^2 \sigma^h}{\mathrm{d} \boldsymbol{x} \ \mathrm{d}(\cos \boldsymbol{\theta})} = \frac{3}{8} \left(1 + \cos^2 \boldsymbol{\theta} \right) \left[\mathcal{F}_T^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\sin^2 \boldsymbol{\theta} \right) \left[\mathcal{F}_L^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right] + \frac{3}{4} \left(\cos \boldsymbol{\theta} \right) \left[\mathcal{F}_A^h(\boldsymbol{x}) \right$$

...contribution to fragmentation function $F^h(x)$

- = $2p/\sqrt{s}$: fractional momentum of particle
- $= \angle$ (incoming particle, outgoing hadron)

$$\frac{\sigma_{T,L}}{\sigma_{\text{tot}}} \equiv \frac{1}{2} \sum_{h} \int dx \ x \cdot \mathcal{F}_{T,L}^{h}(x)$$

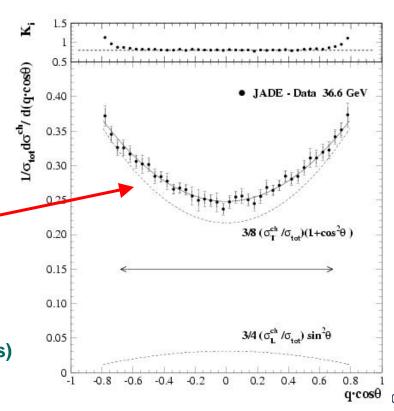
$$\frac{1}{\sigma_{\rm tot}} \cdot \frac{\mathrm{d}\sigma^{\rm ch}}{\mathrm{d}(\mathbf{q} \cdot \cos \theta)} = \frac{3}{8} \eta^{\rm ch} \left[\frac{\sigma_L}{\sigma_{\rm tot}} \left(1 - 3 \cos^2 \theta \right) + \left(1 + \cos^2 \theta \right) \right]$$

- measure cos(q) distribution of charged particles
- fit r_L/r_{tot} and h^{ch} (corrects for neutral particles)

contribution from gluon radiation in quark/antiquark system

asymmetric...

not considered because no experimental distinction between quark/anti-quark



Results

$$\rho_L/\rho_{tot} = 0.067 \pm 0.011$$

Dominant errors:

- limited data statistics (combined 35+44GeV analysis)
- limited MC statistics (preprocessed samples)

$$\left(\frac{\sigma_L}{\sigma_{\text{tot}}}\right)_{\text{PT}} = \frac{\alpha_S}{\pi} + 8.444 \left(\frac{\alpha_S}{\pi}\right)^2$$

$$\alpha_{\text{S}}(36.6 \text{ GeV})\text{=}~0.150 \pm 0.020$$

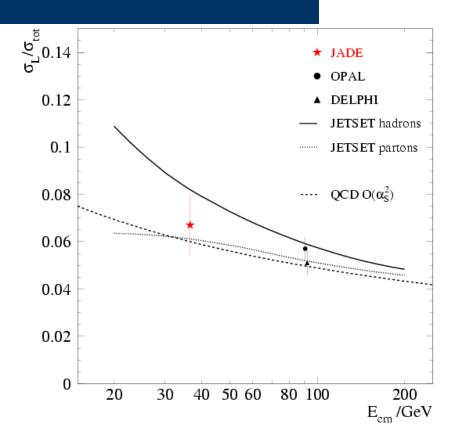
Power corrections:

$$\alpha_{\rm S}({\rm M_Z}) = 0.126 \pm 0.020$$

 $\alpha_{\rm O}(2{\rm GeV}) = 0.3 \pm 0.3$

... not fixed as yet due to low data+MC statistics

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$$\frac{\sigma_L}{\sigma_{\text{tot}}} = \left(\frac{\sigma_L}{\sigma_{\text{tot}}}\right)_{\text{PT}} + a_{\sigma_L} \cdot \frac{16\mathcal{M}}{3\pi^2} \frac{\mu_I}{\sqrt{s}} \cdot \left(\alpha_0(\mu_I) - \alpha_S(\mu) + \mathcal{O}(\alpha_S^2)\right)$$

x Distribution

Momentum spectrum: $\xi = -\ln(x)$

MLLA calculation (Fong, Webber):

$$F_q(\xi, Y) = \frac{N(Y)}{\sigma \sqrt{2\pi}} \cdot \exp\left(\frac{k}{8} - \frac{s\delta}{2} - \frac{(2+k)\delta^2}{4} + \frac{s\delta^3}{6} + \frac{k\delta^4}{24}\right)$$

$$Y \equiv \ln \frac{\sqrt{\varepsilon}}{2\Lambda_{\text{eff}}}$$

$$\delta \equiv \frac{\xi - \langle \xi \rangle}{\sigma}$$

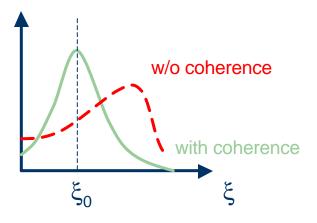
$$\langle \boldsymbol{\xi} \rangle \equiv \langle \boldsymbol{\xi}(\boldsymbol{Y}) \rangle = \frac{\boldsymbol{Y}}{2} \left(1 + \frac{\rho}{24} \sqrt{\frac{48}{\beta \boldsymbol{Y}}} \right) \cdot \left[1 - \frac{\omega}{6 \boldsymbol{Y}} \right] + \mathcal{O}(1)$$

$$\langle \xi_0 - \rangle \langle \xi \rangle \approx \frac{3\rho}{32C_A} \approx 0.35$$

with $N, k, s, \sigma, \beta, \rho, \omega$ known functions of Y, C_A, C_F, N_f

- Test MLLA by fits to measured distributions
 22, 35 and 44 GeV (theory only valid close to ξ₀)
- Free parameters: e.g. N, Λ_{eff} , ξ_0
- Explore the predicted scale dependence of ξ_0

skewed Gaussian



PT prediction

+

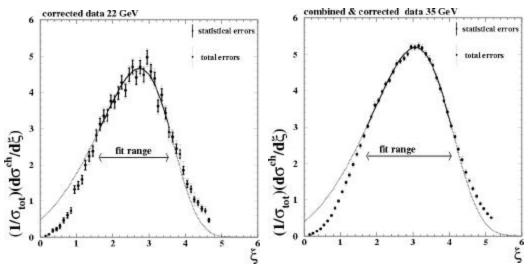
Assume LPHD

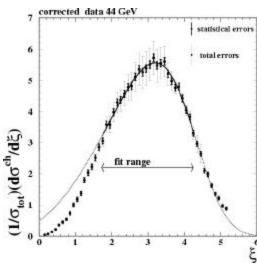
(affects mainly normalisation and not shape)

↓

Hadron spectrum

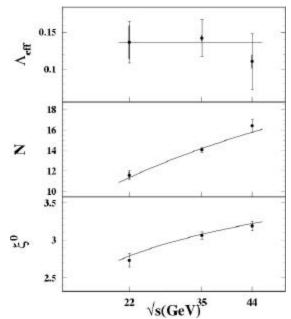
Fits



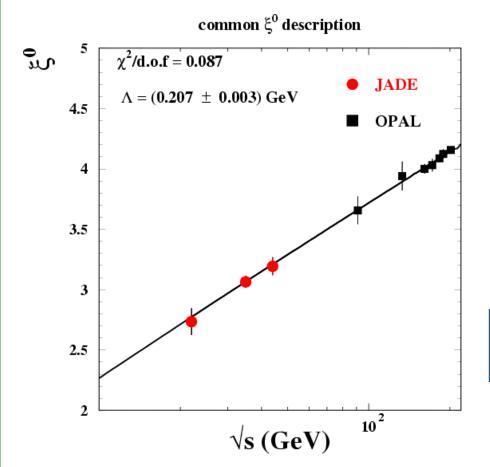


- Good description of data within the kinematic boundaries
- Energy evolution consistent with QCD expectation

	ξ ₀	N	$\Lambda_{ ext{eff}}$
22 GeV	2.74±0.09	11.6±0.4	136±28
35 GeV	3.06±0.05	14.1±0.2	142±25
44 GeV	3.19±0.06	16.4±0.6	110±38



Scale Dependence



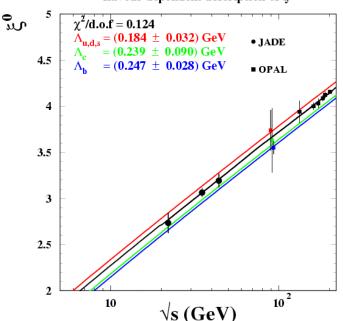
- $\xi_0(Y) = \frac{1}{2}Y + \sqrt{CY} + C$ $Y = \ln(0.5\sqrt{s} / \Lambda_{eff})$
- Use JADE + OPAL data $\sqrt{s} = 22 \dots 202 \text{ GeV}$

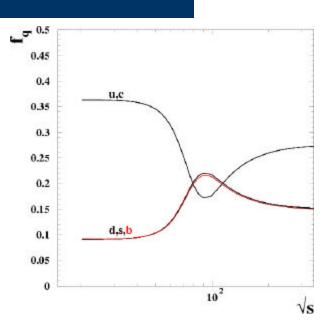
Reasonable description of data: $\Lambda_{\rm eff}$ =207±3 MeV

Flavour dependence

- write ξ_0 (\sqrt{s}) as linear combination of peak positions $\xi_0^{(q)}$ (\sqrt{s}) for flavour q, weighted with branching ratio $f_q(\sqrt{s})$
- $\xi_0^{(c,b)}$ $\xi_0^{(uds)} \propto 0.5$ In $(\Lambda^{(c,b)}/\Lambda^{(uds)})$ \Rightarrow flavour dependence of energy evolution
- fix $\xi_0^{\text{(uds)}}$, $\xi_0^{\text{(c)}}$, $\xi_0^{\text{(b)}}$ with OPAL data @ $\sqrt{\text{s}}$ = M_Z
- fit $\Lambda^{\text{(uds)}}$, $\Lambda^{\text{(c)}}$, $\Lambda^{\text{(b)}}$







Mass effects about 20-30%:

$$\Lambda^{\text{(uds)}} = 184 \pm 32 \text{MeV}$$

$$\Lambda^{(c)} = 239 \pm 90 \text{MeV}$$

$$\Lambda^{(b)} = 247 \pm 28 \text{MeV}$$

Conclusions

Reanalysis of JADE data...

- complements state-of-the-art studies from LEP in the lower energy part of the e⁺e⁻ continuum
- provides stringent tests of perturbative and nonperturbative aspects of QCD
- is needed for constraining (future!) QCD predictions

Keep the data and the software alive since QCD is still in progress!